

RIGID BODY DYNAMICS

Purpose:

- To Study Kinetic States and Principles of Rigid Bodies.

Topics:

- Kinetic States of a Rigid Body.
- Kinetic Principles of a Rigid Body.
- Rigid Body Rotation about an Invariant Axis.



Kinetic States of a Rigid Body:

Theorem-23: The Momentum of a rigid body is equal to its mass times the velocity of the mass center.

$$\underline{P} = m \underline{v}^C \quad \text{or} \quad P_i = m v_i^C \quad (8.1)$$

Theorem-24: The Central Moment of Momentum for a rigid body is equal to its central inertia tensor times the angular velocity vector.

$$\underline{H}^C = \underline{I}^C \cdot \underline{\omega} \quad \text{or} \quad H_i^C = I_{ij}^C \omega_j \quad i, j = 1, 2, 3 \quad (8.2)$$



Equation (8.2) in expanded form is:

$$\begin{aligned} H_1^C &= I_{11}^C \omega_1 + I_{12}^C \omega_2 + I_{13}^C \omega_3 \\ H_2^C &= I_{21}^C \omega_1 + I_{22}^C \omega_2 + I_{23}^C \omega_3 \\ H_3^C &= I_{31}^C \omega_1 + I_{32}^C \omega_2 + I_{33}^C \omega_3 \end{aligned} \quad (8.3)$$

Theorem-25: The Moment of Momentum vector about a general body point “A” in the rigid body can be expressed as:

$$\underline{H}^A = m \underline{\rho}^C \times \underline{v}^A + \underline{I}^A \cdot \underline{\omega} \quad (8.4)$$



Proof:

Consider the rigid body shown;

$\{x_i\}$: its origin is located at the fixed point “O”.

For a rigid body, we have:

$$m = \int_m dm$$

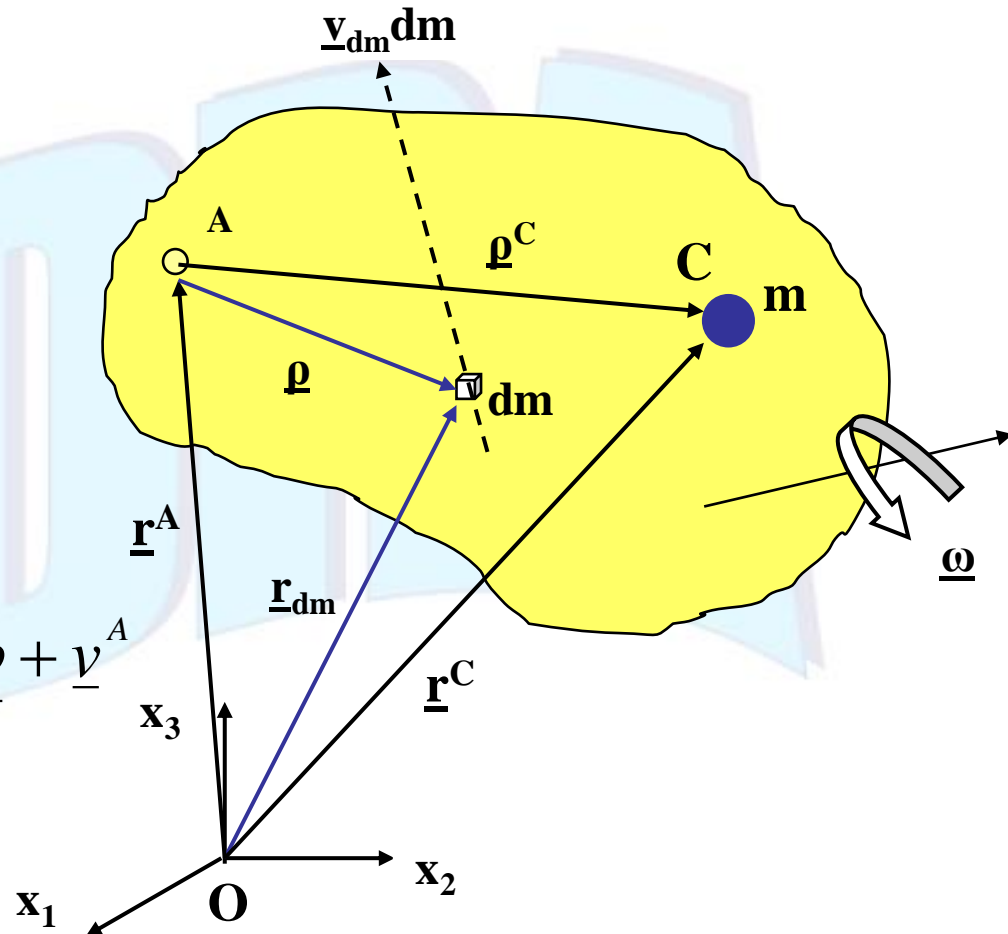
$$\underline{P} = m\underline{v}^C = \int_m \underline{P}_{dm} = \int_m \underline{v}_{dm} dm$$

$$\underline{H}^A = \int_m \underline{H}_{dm}^A = \int_m \underline{\rho} \times \underline{v}_{dm} dm$$

but;

$$\underline{r}_{dm} = \underline{\rho} + \underline{r}^A \quad \text{and} \quad \underline{v}_{dm} = \underline{\dot{\rho}} + \underline{v}^A$$

$$\underline{H}^A = \int_m \underline{\rho} dm \times \underline{v}^A + \int_m \underline{\rho} \times \underline{\dot{\rho}} dm$$



and from the definition of mass center we have:

$$\underline{m}\underline{\rho}^C = \int_m \underline{\rho} dm, \text{ substituting results;}$$

$$\underline{H}^A = \underline{m}\underline{\rho}^C \times \underline{v}^A + \int_m \underline{\rho} \times (\underline{\omega} \times \underline{\rho}) dm \quad (8.5)$$

Since $(\underline{\rho} = \underline{r}_{dm} - \underline{r}^A)$, or $(\rho_i = x_i - x_i^A)$ is a constant magnitude vector, and the point “A” is a body point, then $(\dot{\underline{\rho}} = \underline{\omega} \times \underline{\rho})$ Now, the i-th component of the 2nd term in equation (8.5) may be expressed as:

$$\begin{aligned} \left[\int_m \underline{\rho} \times (\underline{\omega} \times \underline{\rho}) dm \right]_i &= \omega_j \int_m [(x_k - x_k^A)(x_k - x_k^A)\delta_{ij} - (x_i - x_i^A)(x_j - x_j^A)] dm \\ &= I_{ij}^A \omega_j \end{aligned}$$



Substituting the result into equation (8.5), we have:

$$\underline{H}^A = m \underline{\rho}^C \times \underline{v}^A + \underline{I}^A \cdot \underline{\omega} \quad (8.4)$$

However, if point “A” is selected such that

$$(\underline{v}^A = 0), \quad or(A \equiv C), \quad or(\underline{\rho}^C \parallel \underline{v}^A)$$

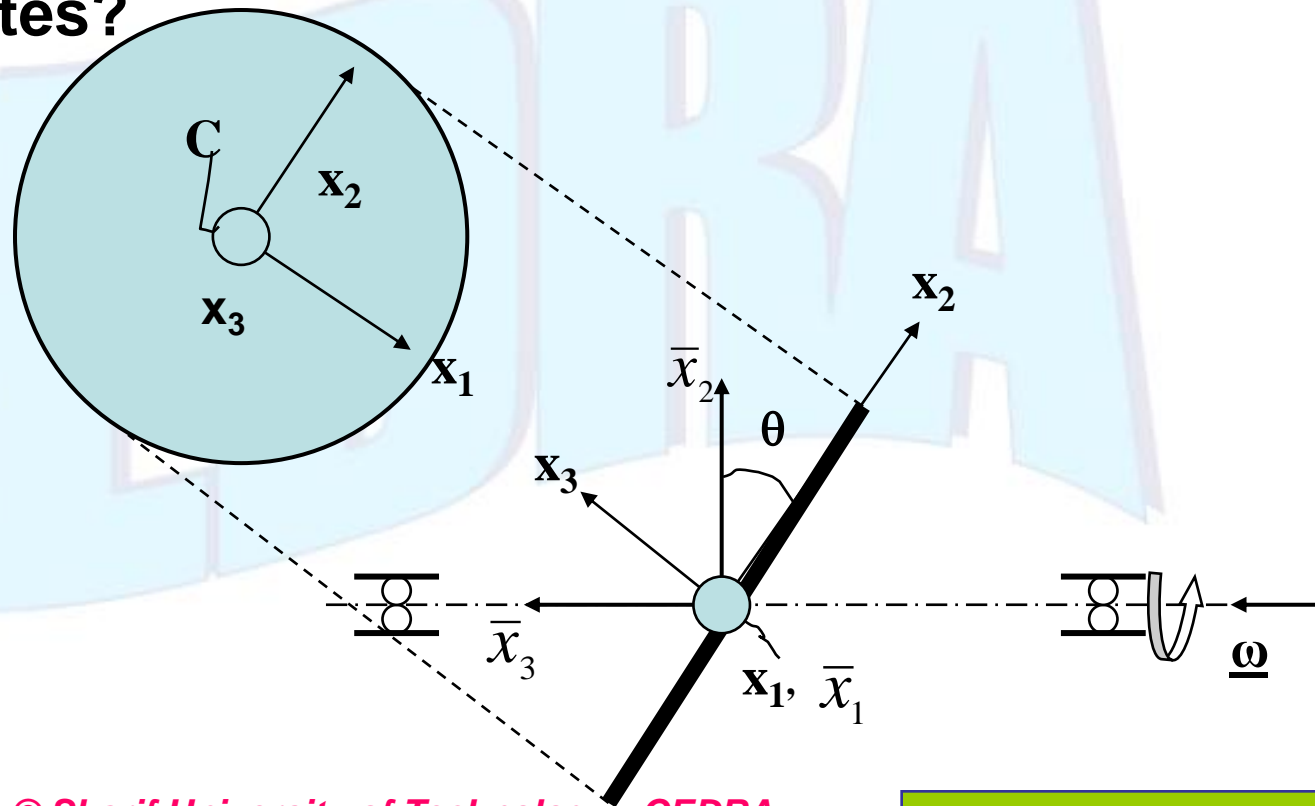
,Then equation (8.4) reduces to:

$$\underline{H}^A = \underline{I}^A \cdot \underline{\omega} \quad or \quad H_i^A = I_{ij}^A \omega_j \quad (8.6)$$



Example: A thin disk of mass “m” and radius “R” is spinning about a fixed axis as shown. Determine the Central Moment of Momentum in terms of:

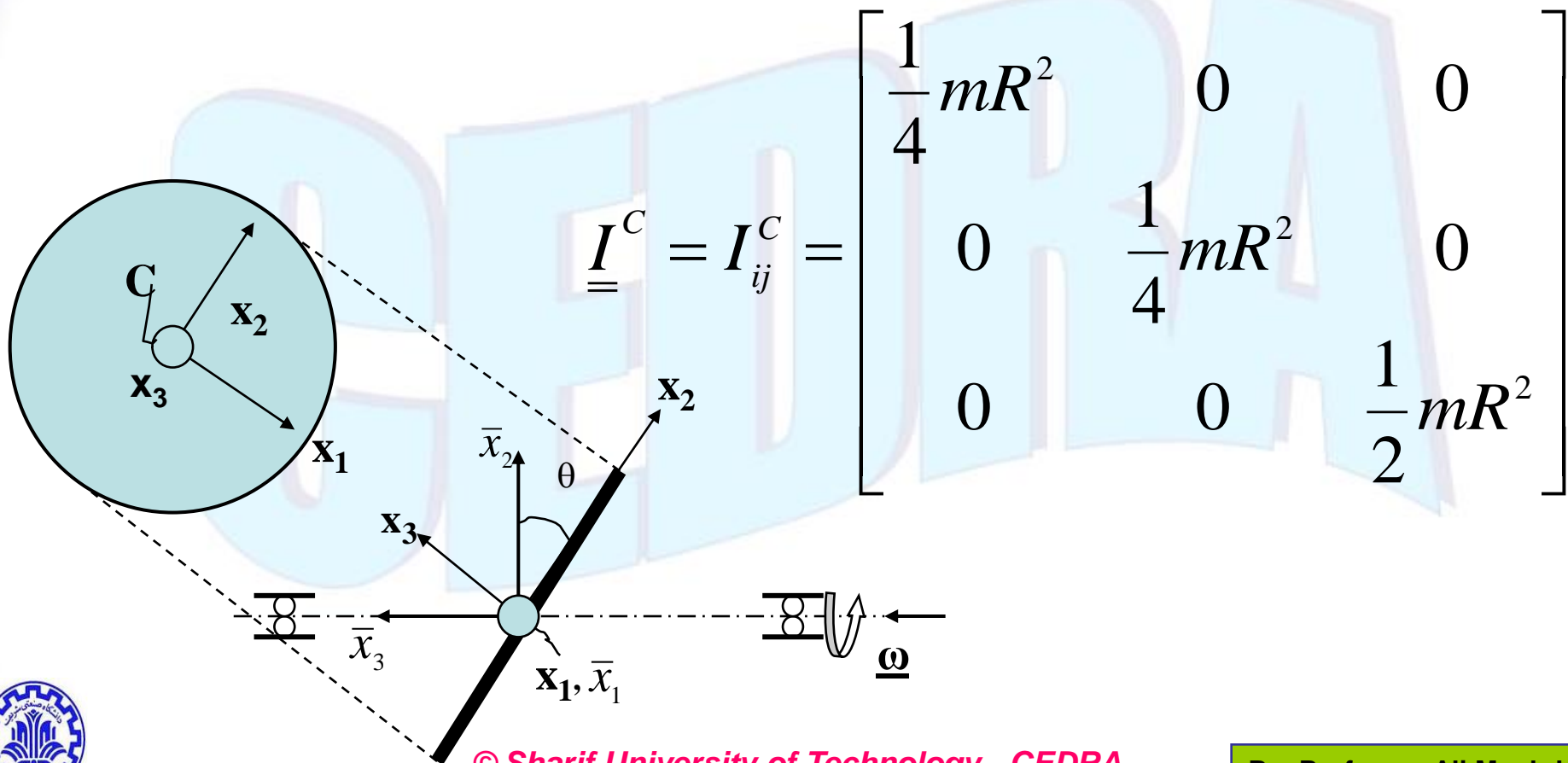
- (a). the principal coordinates? (neglect the disk thickness).
- (b). the coordinate with the rotating axis being as one of the coordinates?



Solution:

(a). In terms of principal coordinates $\{x_i\}$:

From symmetry, the coordinate $\{x_i\}$ is established in the principal directions;



The diagram illustrates a rotating disk of mass m and radius R . A coordinate system $\{\bar{x}_1, \bar{x}_2, \bar{x}_3\}$ is established with the origin at the center of mass C . The \bar{x}_1 and \bar{x}_2 axes are in the plane of the disk, while \bar{x}_3 is along the axis of rotation. The disk is rotating with angular velocity $\underline{\omega}$ about the \bar{x}_3 axis. The principal moments of inertia are given by the matrix $\underline{I}^C = I_{ij}^C$:

$$\underline{I}^C = I_{ij}^C = \begin{bmatrix} \frac{1}{4}mR^2 & 0 & 0 \\ 0 & \frac{1}{4}mR^2 & 0 \\ 0 & 0 & \frac{1}{2}mR^2 \end{bmatrix}$$

The diagram also shows a local coordinate system $\{x_1, x_2, x_3\}$ at the center of the disk, where x_1 and x_2 are in the plane of the disk and x_3 is along the axis of rotation. The angle θ is shown between the \bar{x}_2 axis and the x_2 axis.

The angular velocity $\underline{\omega}$ in terms of principal coordinates is:

$$\{\omega_i\} = \begin{Bmatrix} 0 \\ -\omega \sin \theta \\ \omega \cos \theta \end{Bmatrix}, \text{ then;}$$

$$H_i^c = I_{ij}^c \omega_j = \begin{Bmatrix} 0 \\ -\frac{1}{4} m R^2 \omega \sin \theta \\ \frac{1}{2} m R^2 \omega \cos \theta \end{Bmatrix} \equiv \begin{Bmatrix} H_1^c \\ H_2^c \\ H_3^c \end{Bmatrix} = \underline{H}^c$$



(b). In terms of coordinates oriented along the rotating axis $\{\bar{x}_p\}$; Let us first setup the direction cosines such that:

$$\underline{\underline{T}} = \{\ell_{ip}\} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c\theta & -s\theta \\ 0 & s\theta & c\theta \end{bmatrix}$$

Then, using the transform equation (7.25), we can write:

$$\underline{\underline{\bar{I}}}^c = \underline{\underline{T}}^t \underline{\underline{I}}^c \underline{\underline{T}}, \text{ and}$$

$$\underline{\underline{\bar{I}}}^c = \begin{bmatrix} \frac{1}{4}mR^2 & 0 & 0 \\ 0 & \frac{1}{4}mR^2(1 + \sin^2 \theta) & \frac{1}{4}mR^2 \sin \theta \cos \theta \\ 0 & \frac{1}{4}mR^2 \sin \theta \cos \theta & \frac{1}{4}mR^2(1 + \cos^2 \theta) \end{bmatrix}$$



Then, the **Central Moment of Momentum** is:

$$\underline{\underline{H}}^c = \underline{\underline{I}}^c \cdot \underline{\underline{\omega}} ; \text{ where } \underline{\underline{\omega}} = \begin{Bmatrix} 0 \\ 0 \\ \omega \end{Bmatrix}, \text{ and;}$$

$$\underline{\underline{H}}^c = \begin{Bmatrix} \bar{H}_1^c \\ \bar{H}_2^c \\ \bar{H}_3^c \end{Bmatrix} = \begin{Bmatrix} 0 \\ \frac{1}{4} m R^2 \omega \sin \theta \cos \theta \\ \frac{1}{4} m R^2 \omega (1 + \cos^2 \theta) \end{Bmatrix}$$



Observations:

- Although the spin axis is fixed, the moment of momentum has components along the other axes.
- Even when “ $\underline{\omega}$ ” is constant, while the moment of momentum components may be constant in the body coordinate $\{x_i\}$, the moment of momentum vector will rotate about the rotation axis at a speed of “ ω ”.
- The time rate of change of moment of momentum due to the change in direction is an indication of presence of an exterior moment on the body.



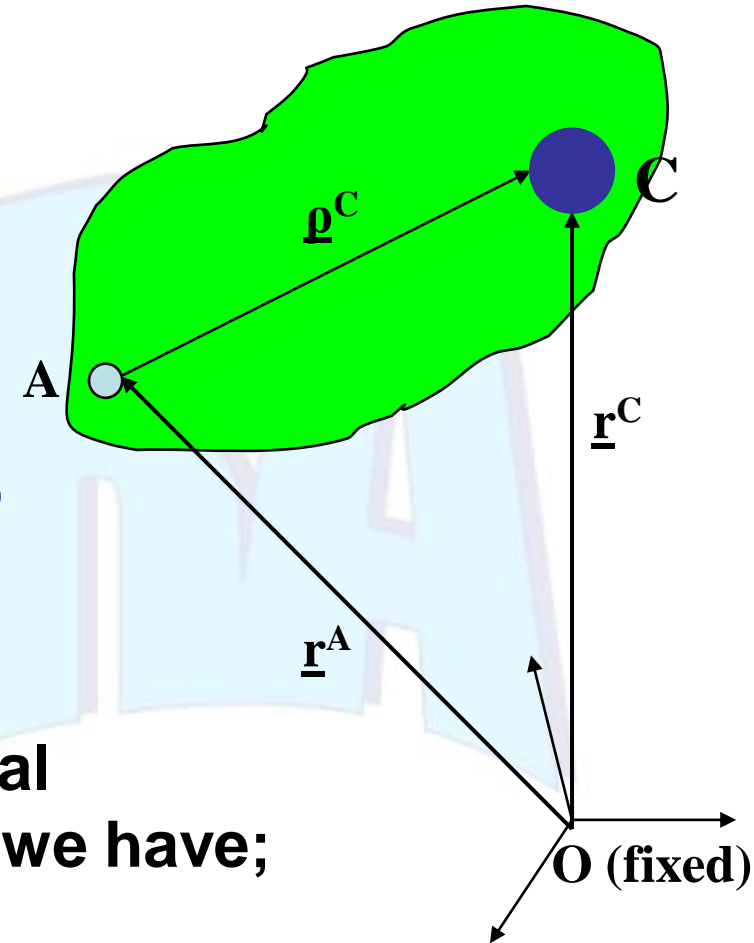
Kinetic Principles of a Rigid Body: Consider a rigid body as shown, where point “A” is a body point, then:

Momentum Principle (**P-Principle**)
for the Rigid Body is:

$$\underline{f} = m \underline{a}^C = \dot{\underline{P}} \quad (8.7)$$

Momentum of Momentum Principle
(**H-Principle**) for the Rigid Body is:

In Chapter-6 we showed that for a system of particles, with the general moment center “A”, and $\underline{P} = m \underline{v}^C$, we have;



$$\underline{M}^A = \dot{\underline{H}}^A + (\underline{v}^A \times \underline{P}) \quad (6.23)$$



However, if

$$\left\{ \begin{array}{l} \underline{v}^A = 0 \quad \text{or} \quad ("A" \text{--}is \text{--}a \text{--}fixed \text{--}point) \\ A \equiv C \\ \underline{v}^A \parallel \underline{P} \end{array} \right\} \Rightarrow \underline{M}^A = \underline{\dot{H}}^A$$

Now for a rigid body, the moment center “A” is also a body point, hence:

$$\begin{aligned} \underline{H}^A &= \underline{\rho}^C \times m \underline{v}^A + \underline{I}^A \cdot \underline{\omega} \equiv (\underline{r}^C - \underline{r}^A) \times m \underline{v}^A + \underline{I}^A \cdot \underline{\omega}, \text{ and;} \\ \underline{\dot{H}}^A &= (\underline{v}^C - \underline{v}^A) \times m \underline{v}^A + (\underline{r}^C - \underline{r}^A) \times m \underline{a}^A + \frac{d}{dt} (\underline{I}^A \cdot \underline{\omega}) \end{aligned}$$

(8.8)



However, if

$$\left\{ \begin{array}{l} \text{"A": is a fixed point, or} \\ A \equiv C, \text{ or} \\ \underline{v}^A = 0 \quad \text{or} \quad \underline{v}^A \parallel \underline{v}^C, \text{ and} \quad \underline{\rho}^C \parallel \underline{a}^A \end{array} \right\} \Rightarrow \underline{M}^A = \underline{\dot{H}}^A = \frac{d}{dt}(\underline{I}^A \cdot \underline{\omega})$$

Then:

$$\underline{M}^A = \underline{\dot{H}}^A = \frac{d}{dt}(\underline{I}^A \cdot \underline{\omega}), \quad \text{or} \quad M_i^A = \frac{d}{dt}(I_{ij}^A \omega_j) \quad (8.9)$$

Equation (8.9) is referred to as the **Euler's Equation**.

Definition: A moment center satisfying the conditions stated above is called an **Eulerian Moment Center**.



Euler's Equation as a function of coordinate variables:

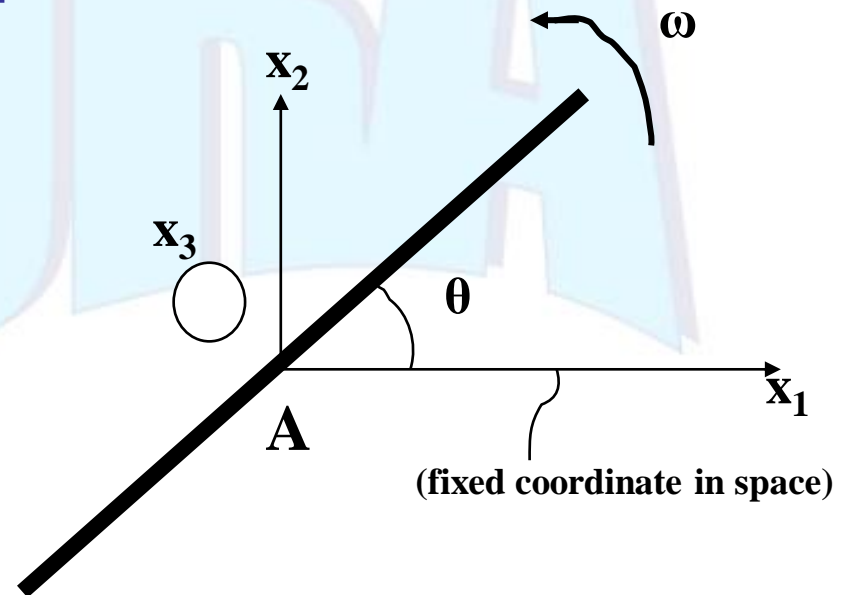
$$M_i^A = \frac{d}{dt} [I_{ij}^A(x) \omega_j]; \quad \text{for } \underline{\text{fixed}} \{x_i\}, \text{ where:} \quad (8.10)$$

$I_{ij}^A(x)$: are components of inertia tensor
(function of coordinate variables)

Example: A Thin Uniform Rod:

$$I_{=\alpha}^A = \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{ml^2}{12} & 0 \\ 0 & 0 & \frac{ml^2}{12} \end{bmatrix}$$

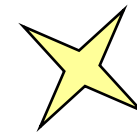
(i.e. $I_{11}^A(\theta), I_{12}^A(\theta), \dots$), where:



$$I_{ij}^A(\theta) = \underline{T}^t \underline{I}_{\alpha}^A \underline{T} = \begin{bmatrix} C\theta & S\theta & 0 \\ -S\theta & C\theta & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{m\ell^2}{12} & 0 \\ 0 & 0 & \frac{m\ell^2}{12} \end{bmatrix} \begin{bmatrix} C\theta & -S\theta & 0 \\ S\theta & C\theta & 0 \\ 0 & 0 & 1 \end{bmatrix} =$$

$$= \begin{bmatrix} \frac{m\ell^2}{12} S^2\theta & \frac{m\ell^2}{12} S\theta C\theta & 0 \\ \frac{m\ell^2}{12} S\theta C\theta & \frac{m\ell^2}{12} C^2\theta & 0 \\ 0 & 0 & \frac{m\ell^2}{12} \end{bmatrix} = \frac{m\ell^2}{12} \begin{bmatrix} S^2\theta & S\theta C\theta & 0 \\ S\theta C\theta & C^2\theta & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Therefore; $I_{11}^A(\theta) = \frac{1}{12} m\ell^2 \sin^2\theta$, and etc.. **Euler's Equation**
 in this form is coordinate dependent.





مفتش كرم