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The Kinetic Principles in the Newtonian Reference Frame (NRF) :

The kinetic state of a material system is conserved unless disturbed by exterior actions. The *Principle of Momentum* (P.M.) and the *Principle of Moment of Momentum* (P.M.M) govern the change.

Recall: Newtonian (Inertial) Reference Frame;
Non-accelerating & Irrotational reference frame.

Admissible Newtonian Forces in the **NRF** shown on a *Free-Body-Diagram* are:

- Contact Forces
- Field Forces (i.e. gravitational field, and electromagnetic field)
- Spring, and Friction Forces



Momentum Principle (M.P.) for a Single Particle: The time rate of change in the momentum of a particle is equal to the resultant external force acting on the particle.

$$\underline{F} = m\underline{a} = \frac{d}{dt}(m\underline{v}) = \underline{\dot{P}} \quad \Rightarrow \quad \underline{F} = \underline{\dot{P}} \quad (6.20)$$

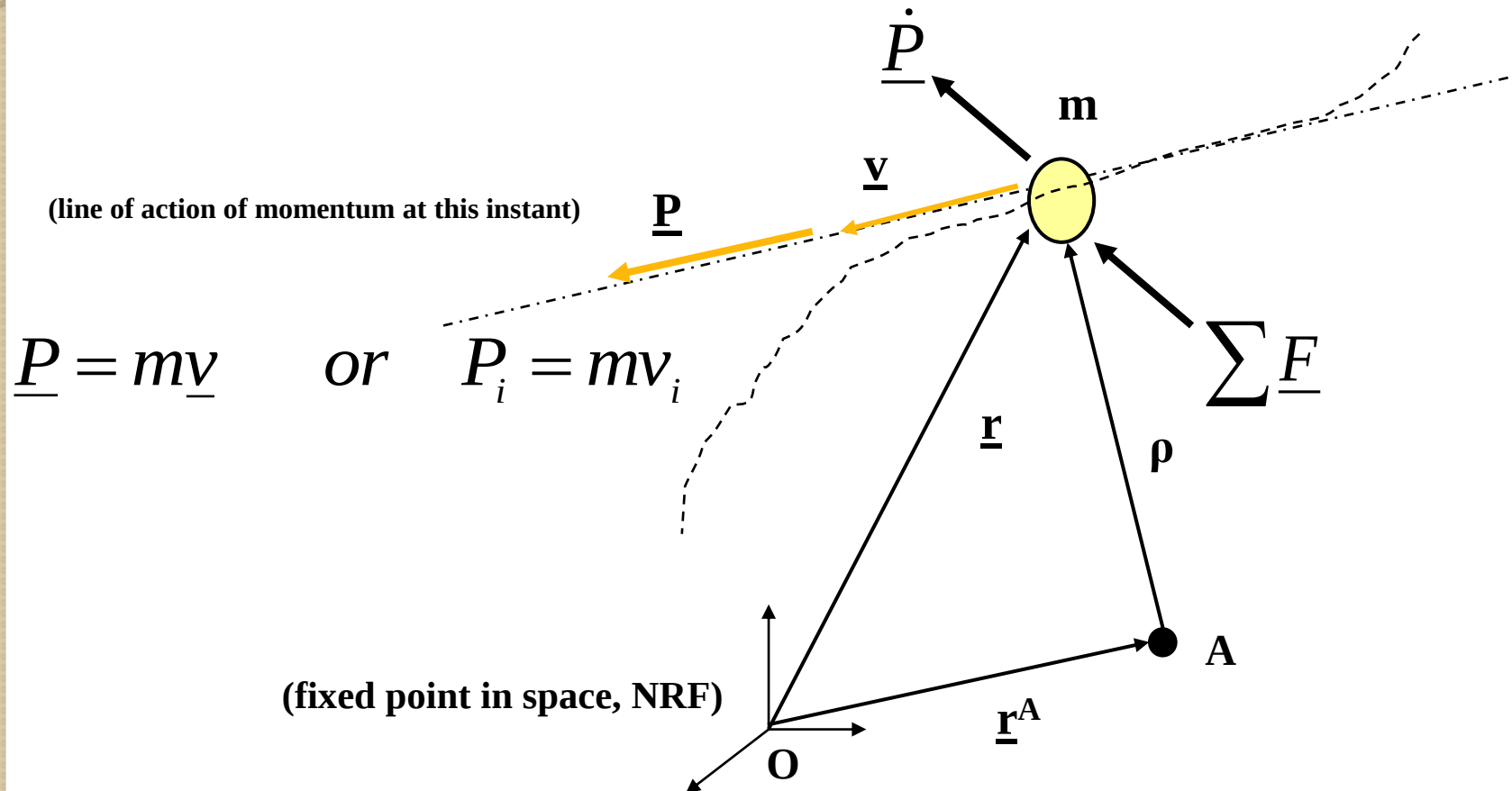
Momentum Principle (M.P.) for a System of Particles:

$$\underline{F} = \sum_{\beta=1}^N \underline{\dot{P}}^{\beta} = \underline{\dot{P}} \quad ; \text{ where } \underline{P} = (\text{Global Momentum})$$



Moment of Momentum Principle (M.M.P.) for a Single Particle:

The time rate of change in moment of momentum of a particle about a fixed moment center “O” in space is equal to the resultant external moment (of forces) about the same moment center.



$$\underline{M}^O = \underline{r} \times \underline{F} = \underline{r} \times \underline{\dot{P}}$$

$$\underline{H}^O = \underline{r} \times \underline{P} = \underline{r} \times m\underline{v}$$

$$\underline{\dot{H}}^O = \frac{d}{dt}(\underline{r} \times \underline{P}) = \underline{\dot{r}} \times \underline{P} + \underline{r} \times \underline{\dot{P}} = \underline{\dot{r}} \times m\underline{\dot{r}} + \underline{M}^O$$

0

$$\underline{M}^O = \underline{\dot{H}}^O \quad (6.21)$$

Remark: The Moment of Momentum Principle can actually be expressed more generally for any **moment center “A”** which has a **zero velocity** in space.

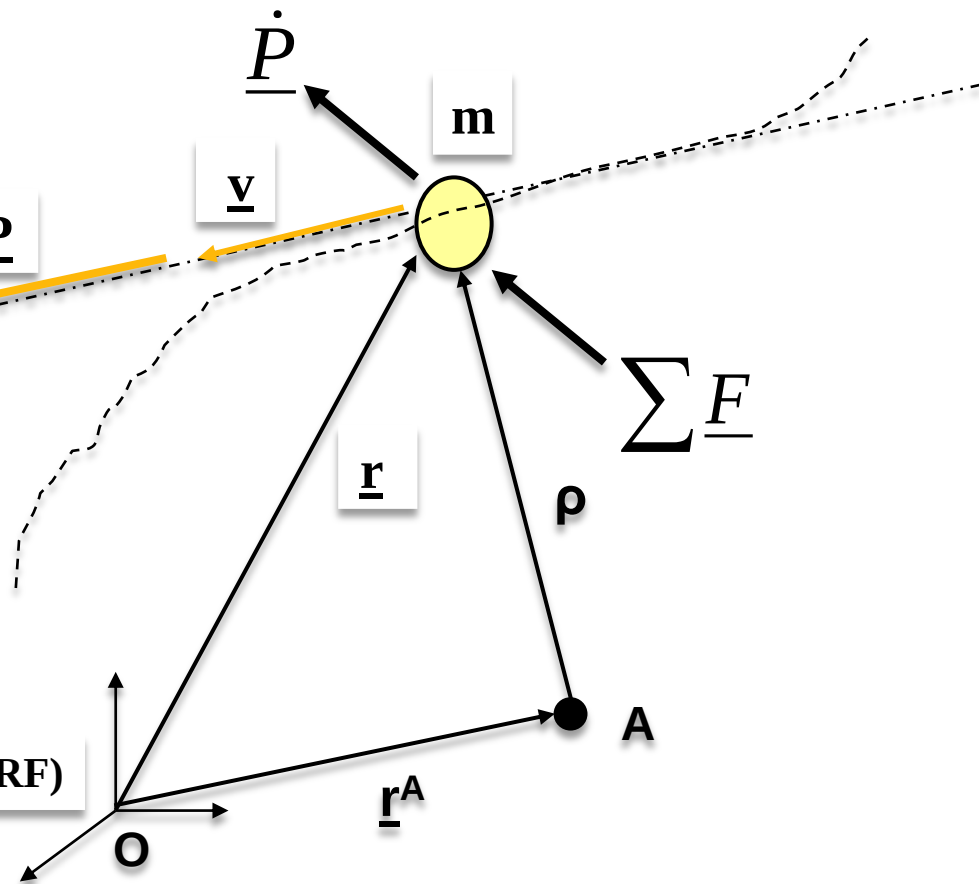
$$\underline{M}^A = \underline{\dot{H}}^A; \text{A} = \text{moment center with zero velocity, but may be accelerating.}$$



(line of action of momentum at this instant)

$$\underline{P} = m\underline{v} \quad \text{or} \quad P_i = mv_i$$

(fixed point in space, NRF)



Proof:

$$\begin{aligned}\underline{H}^A &= \underline{\rho} \times \underline{P} = (\underline{r} - \underline{r}^A) \times \underline{P} \\ \dot{\underline{H}}^A &= \dot{\underline{r}} \times \underline{P} - \dot{\underline{r}}^A \times \underline{P} + (\underline{r} - \underline{r}^A) \times \dot{\underline{P}} = \\ &= -\underline{v}^A \times \underline{P} + \underline{\rho} \times \underline{F} = \\ &= -\underline{v}^A \times \underline{P} + \underline{M}^A \\ \dot{\underline{H}}^A &= \underline{M}^A - \underline{v}^A \times \underline{P}\end{aligned}$$



If $\underline{v}^A = 0$, or $\underline{v}^A \parallel \underline{P} \Rightarrow \dot{\underline{H}}^A = \underline{M}^A$

(A is an admissible moment center)

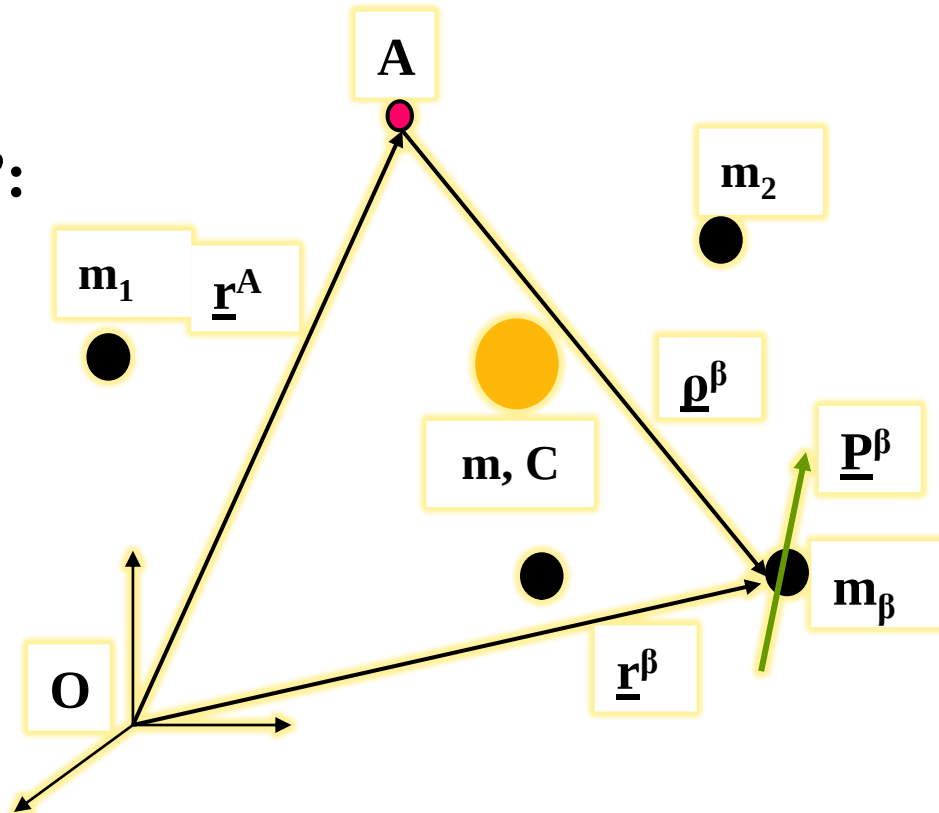


Moment of Momentum Principle (M.M.P.) for a System of Particles:

➤ About a fixed point “O”:

$$\underline{M}^O = \underline{\dot{H}}^O$$

where: $\left\{ \begin{array}{l} \underline{H}^O = \sum_{\beta=1}^N \underline{r}^{\beta} \times \underline{P}^{\beta}, \\ \text{and : } \underline{M}^O : \text{Resultant Moment} \end{array} \right.$



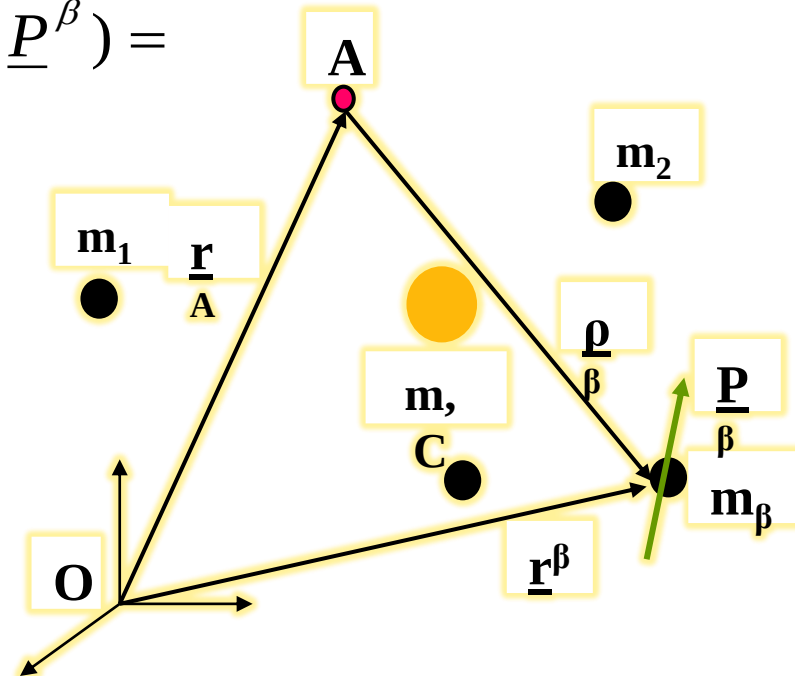
➤ About a general point “A”:

$$\underline{\dot{H}}^A = \frac{d}{dt} \sum_{\beta=1}^N \underline{\rho}^\beta \times \underline{P}^\beta = \sum_{\beta=1}^N \underline{\rho}^\beta \times \underline{\dot{P}}^\beta + \sum_{\beta=1}^N \underline{\dot{\rho}}^\beta \times \underline{P}^\beta =$$

$$= \underline{M}^A + \left(\sum_{\beta=1}^N \underline{\dot{\rho}}^\beta \times \underline{P}^\beta - \sum_{\beta=1}^N \underline{\dot{r}}^A \times \underline{P}^\beta \right) =$$

$$= \underline{M}^A - (\underline{v}^A \times \sum_{\beta=1}^N \underline{P}^\beta)$$

$$\underline{\dot{H}}^A = \underline{M}^A - (\underline{v}^A \times \underline{P})$$



Since: $\underline{\rho}^\beta = \underline{r}^\beta - \underline{r}^A \Rightarrow \underline{\dot{\rho}}^\beta = \underline{\dot{r}}^\beta - \underline{\dot{r}}^A$



$$\underline{\dot{H}}^A = \underline{M}^A - (\underline{v}^A \times \underline{P}) \quad (6.23)$$

However, if

$$\left\{ \begin{array}{l} \underline{v}^A = 0 \\ A \equiv C \\ \underline{v}^A \parallel \underline{P} \end{array} \right\} \Rightarrow \underline{\dot{H}}^A = \underline{M}^A \quad (6.24)$$

Since for a constant mass system: $\underline{P} = m\underline{v}^C$ Point “A” is an admissible moment center for equation (6.24), if and only if it satisfies one of the above conditions.



Linear and Angular Impulse:

- Linear Impulse of a Force “F”: Let; $\underline{F} = \underline{F}(t)$ and $\underline{P} = m\underline{v}$, then:

$$\underline{F} = \dot{\underline{P}} = \frac{d}{dt} \underline{P} \Rightarrow \underline{F} dt = d \underline{P} \Rightarrow$$

$$\underline{f} = \underline{\text{impulsive} - \text{force}} = \int_{t_1}^{t_2} \underline{F} dt = \int_{t_1}^{t_2} d \underline{P} = \underline{P}(t_2) - \underline{P}(t_1) \quad (6.25)$$

- Angular Impulse of a Moment “M”: Let; $\underline{M}^O = \underline{M}^O(t)$ and $\underline{H}^O = \underline{r} \times \underline{P}$, then:

$$\underline{M}^O = \dot{\underline{H}}^O = \frac{d}{dt} \underline{H}^O \Rightarrow \underline{M}^O dt = d \underline{H}^O \Rightarrow \quad (6.26)$$

$$\underline{M} = \underline{\text{impulsive} - \text{moment}} = \int_{t_1}^{t_2} \underline{M}^O dt = \int_{t_1}^{t_2} d \underline{H}^O = \underline{H}^O(t_2) - \underline{H}^O(t_1)$$



Note: *the “**Momentum**” and the “**Moment of Momentum**”*
*are conserved if the **Impulse Integral Vanishes**, that is when:*

1. $\underline{F} = \underline{M}^0 = 0$ (the integrand vanishes).
2. The integrand is finite, but the time interval is infinitesimal ($\Delta t = \varepsilon$), (i.e. impact problems, and explosions).
3. The integrand is a cyclic function and the time interval is one complete period, (i.e. materials under cyclic loading).



Ex: A simple pendulum constrained to oscillate in a plane as shown, has a mass “m” and a length “L”. Derive its differential equation of motion?

Solution:

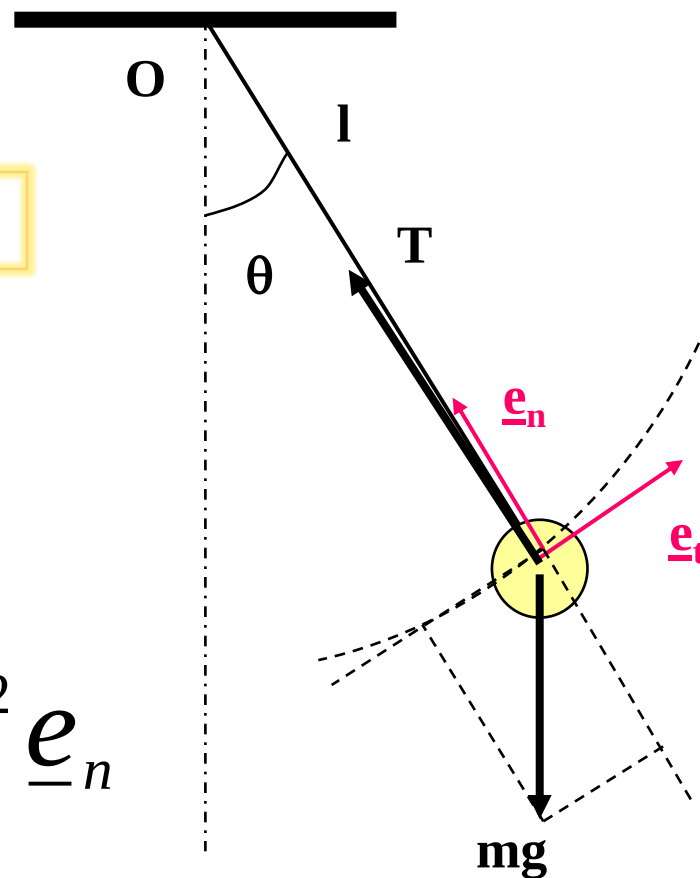
(a). By the **Momentum Principle:**

❖ **Motion Analysis:**

$$\underline{v} = \ell \dot{\theta} \underline{e}_t$$

$$\underline{a} = \ell \ddot{\theta} \underline{e}_t + \ell \dot{\theta}^2 \underline{e}_n$$

❖ **Free Body Diagram**



❖ Kinetics (M.P.):

$$\underline{f} = \underline{\dot{P}} = m\underline{a}$$

$$T\underline{e}_n - mg(\cos\theta\underline{e}_n + \sin\theta\underline{e}_t) = m\underline{a} = m(\ell\ddot{\theta}\underline{e}_t + \ell\dot{\theta}^2\underline{e}_n)$$

$$\left\{ \begin{array}{l} \underline{e}_t : \quad -mg \sin \theta = m\ell \ddot{\theta} \quad \Rightarrow \ddot{\theta} + \frac{g}{\ell} \sin \theta = 0 \quad \leftarrow \underline{\text{D.E.M.}} \\ \underline{e}_n : \quad T - mg \cos \theta = m\ell \dot{\theta}^2 \quad \Rightarrow T = m(\ell \dot{\theta}^2 + g \cos \theta) \quad \leftarrow \end{array} \right.$$



(b). By the Momentum of Momentum Principle:

❖ Motion Analysis:

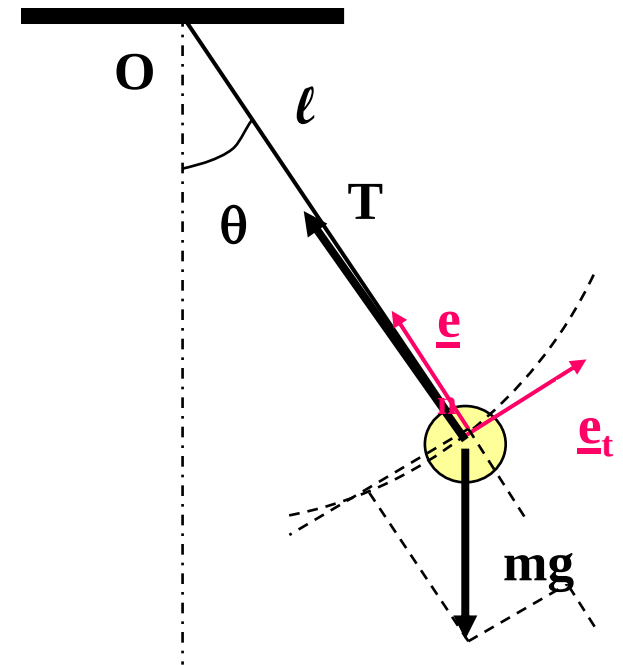
$$\underline{v} = \ell \dot{\theta} \underline{e}_t$$

❖ Free Body Diagram

❖ Kinetics (M.M.P.):

$$\left\{ \begin{array}{l} \underline{P} = m\underline{v} = m\ell \dot{\theta} \underline{e}_t \\ \underline{H}^O = \underline{r} \times \underline{P} = (-\ell \underline{e}_n) \times (m\ell \dot{\theta} \underline{e}_t) = m\ell^2 \dot{\theta} \underline{e}_b \end{array} \right.$$

$$\left\{ \begin{array}{l} \underline{M}^O = \underline{\dot{H}}^O \Rightarrow -mg\ell \sin \theta = m\ell^2 \ddot{\theta} \\ \ddot{\theta} + \frac{g}{\ell} \sin \theta = 0 \end{array} \right.$$



Ex: A horizontal table is spinning at a constant angular velocity of “ Ω ”. On the table, a spring-retained mass “ m ” is oscillating in a straight slot as shown. The spring constant of each spring is “ k ”. Determine the equation of motion of the mass, if the coefficient of friction between the mass and the slot is “ μ ”.

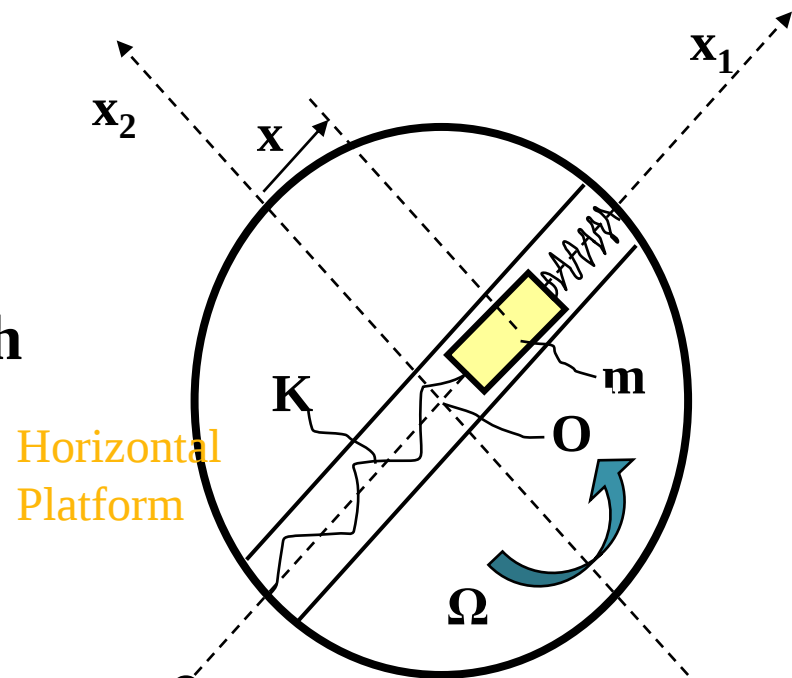
Solution:

❖ **Motion Analysis:**

a). Fix KRF $\{x_i\}$ to the table with reference point at “O”.

b). KRF Motion;

$$\underline{\Omega} = \Omega \underline{e}_3, \quad \dot{\underline{\Omega}} = 0, \quad \underline{v}_O = \underline{a}_O = 0$$



c). Relative Motion;

$$\underline{\bar{r}} = x \underline{e}_1, \quad \underline{\bar{v}} = \dot{x} \underline{e}_1, \quad \underline{\bar{a}} = \ddot{x} \underline{e}_1$$

d). Absolute Motion;

$$\underline{a}_m = \underline{a}_O + \underline{\bar{a}} + \underline{\dot{\Omega}} \times \underline{\bar{r}} + 2\underline{\Omega} \times \underline{\bar{v}} + \underline{\Omega} \times (\underline{\Omega} \times \underline{\bar{r}})$$

$$\underline{a}_m = (\ddot{x} - \Omega^2 x) \underline{e}_1 + 2\Omega \dot{x} \underline{e}_2$$

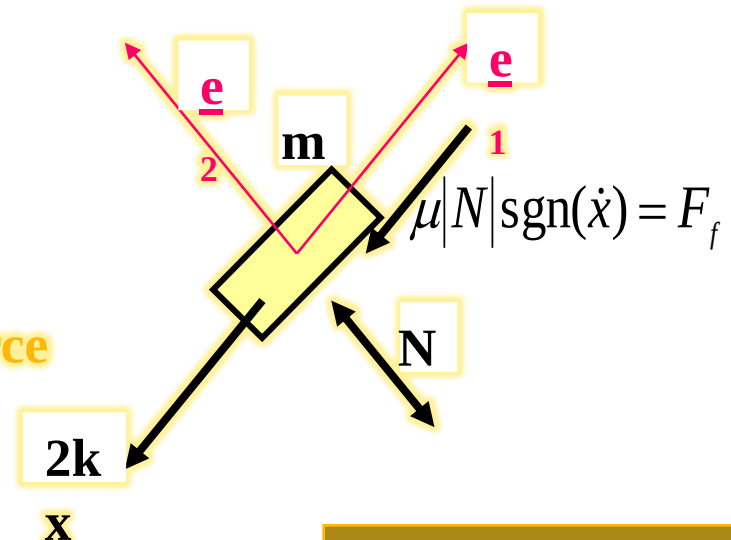


❖ **Free Body Diagram**

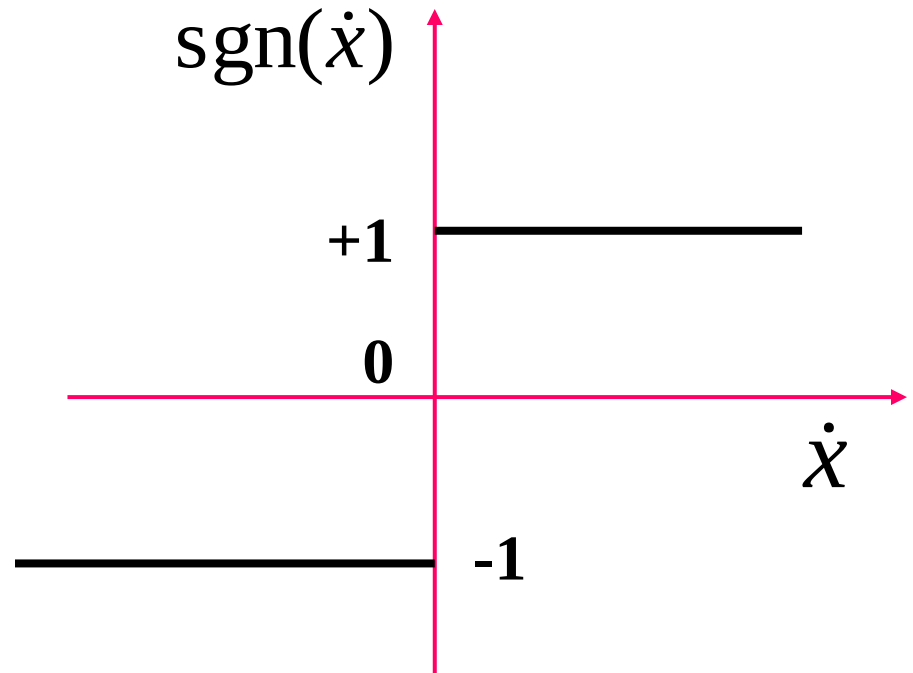
Coulomb (Dry) Friction:

$$\mu |N| \operatorname{sgn}(\dot{x}) = F_f, \text{ a constant force}$$

to oppose the sliding motion.



Signum Function:

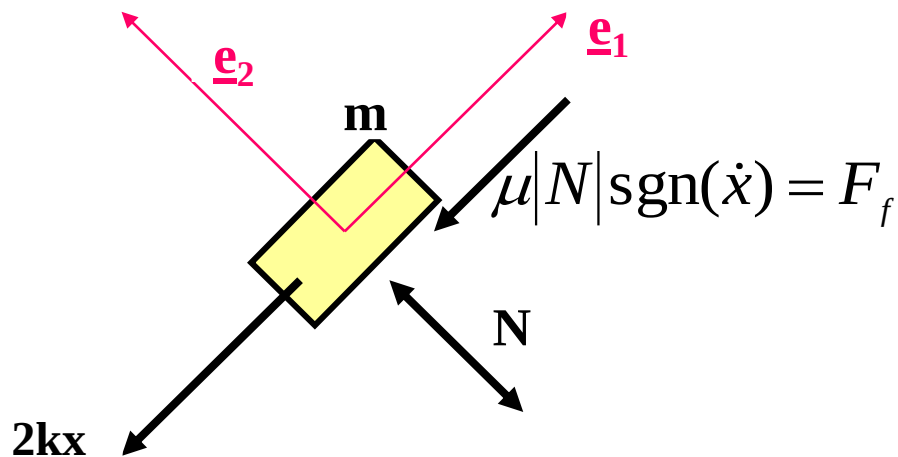


$$\text{sgn}(\dot{x}) = +1 \quad \text{if} \quad \dot{x} > 0$$

$$\text{sgn}(\dot{x}) = -1 \quad \text{if} \quad \dot{x} < 0$$

$$-1 \leq \text{sgn}(\dot{x}) \leq +1 \quad \text{if} \quad \dot{x} = 0$$





❖ Kinetic Analysis

Momentum Principle; $\underline{f} = \underline{\dot{P}} = m\underline{a}$

$$(-2kx - \mu|N|\text{sgn}(\dot{x}))\underline{e}_1 + N\underline{e}_2 = m[(\ddot{x} - \Omega^2 x)\underline{e}_1 + 2\Omega\dot{x}\underline{e}_2]$$

$$\underline{e}_2 : \quad N = 2m\Omega\dot{x}$$

$$\underline{e}_1 : \quad -2kx - \mu|N|\text{sgn}(\dot{x}) = m(\ddot{x} - \Omega^2 x)$$

$$-2kx - 2\mu m|\Omega||\dot{x}|\text{sgn}(\dot{x}) = m(\ddot{x} - \Omega^2 x)$$



$$m\ddot{x} + 2\mu m|\Omega|\dot{x} + (2k - m\Omega^2)x = 0, \quad \text{or}$$

$$\ddot{x} + 2\mu|\Omega|\dot{x} + \left(\frac{2k}{m} - \Omega^2\right)x = 0 \quad \longleftarrow \text{D.E.M.}$$

(Linearly decaying oscillation due to the dry friction). Note that direction of “N” is arbitrary and does not affect the D.E.M.



Moment of Momentum Principle (M.M.P.) in terms of a rotating coordinate (KRF) with angular velocity “ $\underline{\Omega}$ ”:

$$\underline{M}^O = \dot{H}_i^O \underline{u}_i + \underline{\Omega} \times \underline{H}^O \quad \text{or} \quad M_i^O = \dot{H}_i^O + \gamma_{ijk} \Omega_j H_k^O$$

(6.27)



مقدمه