

## Sharif University of Technology School of Mechanical Engineering Center of Excellence in Energy Conversion

## **Advanced Thermodynamics**

Lecture 25

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- **Ø** Ideal solutions:  $\overline{V}_A \overline{V}_A = 0$ , at *P*, *T* and same phase as the mixture.
- Ø Ideal gas is a special case of ideal solutions.

$$\text{ISM} \to V = \overline{V}_A n_A + \overline{V}_B n_B = \overline{v}_A n_A + \overline{v}_B n_B$$

Ø It is similar to Amagat model (but  $\overline{V}_A - \overline{v}_A = 0$  at *P*, *T* and same phase as the mixture!)

 $m_i$ : electrochemical potential of ith component = driving force which causes mass transfer= partial molal Gibbs function

$$\boldsymbol{m}_{i} = \left(\frac{\partial U}{\partial N_{i}}\right)_{S, V, N_{j\neq i}} = \left(\frac{\partial H}{\partial N_{i}}\right)_{S, P, N_{j\neq i}} = \left(\frac{\partial A}{\partial N_{i}}\right)_{T, V, N_{j\neq i}} = \left(\frac{\partial G}{\partial N_{i}}\right)_{T, P, N_{j\neq i}} = \left(\frac{\partial$$

**Ø** For a pure substance (T= const.)

$$d\overline{g}_A = \overline{v}_A dP = \overline{R}T d \ln f_A \text{ and } \lim_{P \to 0} \left(\frac{f_A}{P}\right) = 1$$

 $\boldsymbol{\emptyset}$  For a component in a mixture (T= const.)

$$d\overline{G}_A = \overline{V}_A dP = \overline{R}T d \ln \overline{f}_A \text{ and } \lim_{P \to 0} \left(\frac{\overline{f}_A}{y_A P}\right) = 1$$

 $\overline{f_A}$  = The fugacity of a component at the mixture conditions  $f_A$  = The fugacity of a pure A,  $\overline{f_A} = y_A f_A$ 

$$\Rightarrow \overline{R}T \ d \int_{\ln\left(\overline{f_A}/f_A\right)^* = \ln y_A}^{\ln\left(\overline{f_A}/f_A\right)} \ln\left(\frac{\overline{f_A}}{f_A}\right) \ln = \int_{P^* \to 0}^{P} (\overline{V_A} - \overline{V_A}) dP$$

T and composition= const.  $\rightarrow \overline{R}T \ln\left(\frac{\overline{f_A}}{y_A f_A}\right) = \int_0^P (\overline{V_A} - \overline{V_A}) dP$ 

**Ø** There is a mixture containing 75%  $CH_4$  and 25%  $C_2H_4$ 

at  $T = 25^{\circ}C$  & P = 8.25 MPa. Fin gas mass

Ø I.G. assumption:

$$\begin{aligned} PV &= mRT \text{, using Dalton model} \rightarrow \\ P_{CH_4} &= y_{CH_4}P = 0.75 \times 8.25 = 6.188MPa \& P_{C_2H_4} = 2.063MPa \\ m_{CH_4} &= \frac{P_{CH_4}V}{R_{CH_4}T} = 20.02Kg \& m_{C_2H_4} = \frac{P_{C_2H_4}V}{R_{C_2H_4}T} = 11.07Kg \\ \rightarrow m &= m_{CH_4} + m_{C_2H_4} = 31.09Kg \end{aligned}$$

Ø Using I.S.M. assumption:

$$V = \overline{V_A} n_A + \overline{V_B} n_B \rightarrow \frac{V}{n_t} = (y\overline{v})_{CH_4} + (y\overline{v})_{C_2H_4}$$
$$\overline{V_A} - \overline{V_A} = 0$$

## **Ad. Thermodynamics**

Ø Using I.S.M. assumption:

no chart for  $\overline{v_{CH_4}}$  and  $\overline{v_{C_2H_4}} \rightarrow$  use the generalized charts  $\rightarrow$ 

$$T_{r,CH_4} = \frac{T}{T_{CH_4}} = 1.56, \ P_{r,CH_4} = \frac{8.25}{4.04} = 1.78, \ T_{r_{,C_2H_4}} = 1.06, \ P_{r_{,C_2H_4}} = 1.61$$

Generalized compressibility factor chart  $\rightarrow Z_{CH_4} = 0.88 \& Z_{C_2H_4} = 0.35$ 

$$\overline{v}_{CH_4} = \frac{Z\overline{R}T}{P} = \frac{0.88 \times 8.31434 \times 298.2}{8250} = 0.26457 \ \frac{m^3}{Kmole}$$

$$\overline{v}_{C_{2H_4}} = \frac{Z\overline{R}T}{P} = \frac{0.35 \times 8.31434 \times 298.2}{8250} = 0.1052 \ \frac{m^3}{Kmole}$$

$$\rightarrow \overline{v} = 0.75(0.26457) + 0.25(0.1052) = 0.2247 \ \frac{m^3}{Kmole}$$

$$M = \sum y_i M_i = 0.75 \times 16 + 0.25 \times 28 = 19$$

$$\rightarrow v = \frac{\overline{v}}{M} = 0.011825 \ \frac{m^3}{Kg} \rightarrow m = \frac{V}{v} = \frac{0.5}{0.011825} = 42.3Kg$$

- Ø Example: Dilute non-electrolyte liquid solutions:
  - Ø Water as solvent and salt as a solute (low volume fraction)Index 1 is for solvent and 2 is for solute
- Ø Henry's Law for dilute liquid solution:
- **Ø** This law was presented for partial pressure instead of fugacity at first,  $P_2 = K_2 x_2$ .

- Ø According to Henry's Law :
- $\boldsymbol{\emptyset}$  Gibbs-Duhem relation for const. *P* and *T*

$$\rightarrow \sum_{i=1}^{N} n_i \ d \mathbf{m}_i = 0$$

**Ø** Since  $(d\overline{G}_i)_T = (d \mathbf{m}_i)_T = \overline{R}T d \ln f_i$ 

$$\rightarrow n_1 \overline{R} T \, d \, \ln f_1 + n_2 \overline{R} T \, d \, \ln f_2 = 0$$

$$\rightarrow x_1 \left[ \frac{\partial \left( \ln f_1 \right)}{\partial x_2} \right]_{T,P} + x_2 \left[ \frac{\partial \left( \ln f_2 \right)}{\partial x_2} \right]_{T,P} = 0$$

$$x_1 = 1 \quad \left[ \frac{\partial \ln f_1}{\partial \ln x_1} \right]_{T,P} = \left[ \frac{\partial \ln f_2}{\partial \ln x_2} \right]_{T,P} = 0$$

 $n_1 d \mathbf{m}_1 + n_2 d \mathbf{m}_2 = 0$ 

$$f_{2} = K_{2}x_{2} \rightarrow \ln f_{2} = \ln K_{2} + \ln x_{2} \rightarrow \left[\frac{\partial \ln f_{2}}{\partial \ln x_{2}}\right]_{T,P} = 1$$

$$\left[\frac{\partial \ln f_{1}}{\partial \ln x_{1}}\right]_{T,P} = \left[\frac{\partial \ln f_{2}}{\partial \ln x_{2}}\right]_{T,P} \Rightarrow \left[\frac{\partial \ln f_{1}}{\partial \ln x_{1}}\right]_{T,P} = 1$$

$$\rightarrow \int_{f_{1}^{0}}^{f_{1}} d \ln f_{1} = \int_{x_{1}=1}^{x_{1}} d \ln x_{1} \rightarrow \boxed{f_{1} = f_{1}^{0}x_{1}}$$

$$f_{1}^{0} = \text{Eugacity of pure 1 at the same T and P of solutions.}$$

 $f_1^0$  = Fugacity of pure 1 at the same T and P of solution

- **Ø** If the vapor in equilibrium with the liquid is assumed to be ideal gas  $P_1 = P_1^0 x_1$
- Ø This law, which is a resultant of Henry's law, is named Rault's law for the solvent.
- Ø It may be shown that if the Henry's law is valid for a solute in a dilute solution, the Raul law is valid for the solvent of that solution.