

Sharif University of Technology School of Mechanical Engineering Center of Excellence in Energy Conversion

Advanced Thermodynamics

Lecture 16

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2011

Ad. Thermodynamics

- Ø Microstate: Exact specification of the system at microscopic level, it tells us all needed information.
- Ø Microstate can be expanded by degeneracy or quantum states or molecular way.
- Ø Macrostate: Listing the particles in energy levels for, N_j , with constraints $N = \sum_j N_j$ and $U = \sum_j N_j e_j$

N = No. of particles, U = Energy of the system, and $e_i = Energy$ level

- Ø Thus, a degenerate state is one in which there is more than one way of obtaining an energy level.
- Ø In a non-degenerate state, there is only one way.

- **Ø** Information & constraints: $N = \sum_{j} N_{j}$ and $U = \sum_{j} N_{j} e_{j}$
- Ø Consider a system with two independent particles and three energy levels , e_0 , e_1 , and e_2 with N = 2 and U = 2

Ø For $g_j = 0$:

- Ø Macro I: $N_0 = 1, N_1 = 0, \text{ and } N_2 = 1$
- Ø Macro II: $N_0 = 0, N_1 = 2, \text{ and } N_2 = 0$

- Ø Degeneracy of a system with only translational kinetic energy:
- $\boldsymbol{\emptyset}$ C is the number of the translational kinetic energy.

$$g = \frac{(c+2)(c+1)}{2}$$

- Ø Statistical models
 - Ø Classical mechanics: particles are distinguishable, no limit for
 No. of particles per quantum state, Boltzmann-Model
 - Ø Quantum mechanics: particles are indistinguishable
 - Ø Maximum 1 per quantum state, Paoli (Fermi-Dirac)
 - Ø No limit, Bose-Einstain

w = Thermodynamic Probability

$$W_{Bol.} = N ! \prod_{j} \left[\frac{g_{j}^{N_{j}}}{N_{j}!} \right]$$
$$W_{B.E.} = \prod_{j} \left[\frac{(g_{j} + N_{j} - 1)!}{(g_{j} - 1)!N_{j}!} \right]$$
$$W_{F.D.} = \prod_{j} \left[\frac{g_{j}!}{(g_{j} - N_{j})!N_{j}!} \right]$$
$$W_{Corrected Boltzmann} = \frac{W_{Boltzmann}}{N!} = \prod_{j} \left[\frac{g_{j}^{N_{j}}}{N_{j}!} \right]$$

 \emptyset Note! In reals ystems $N_j \ll g_j$